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Quarter-wave layers with 50% reflectance for obliquely incident unpolarized light

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The conditions under which light interference in a transparent quarter-wave layer of refractive index n_1 on a transparent substrate of refractive index n_2 leads to 50% reflectance for incident unpolarized light at an angle φ are determined. Two distinct solution branches are obtained that correspond to light reflection above and below the polarizing angle, φ_p , of zero reflection for p polarization. The real p and s amplitude reflection coefficients have the same (negative) sign for the solution branch $\varphi > \varphi_p$ and have opposite signs for the solution branch $\varphi < \varphi_p$. Operation at $\varphi < \varphi_p$ is the basis of a 50%–50% beam splitter that divides an incident totally polarized light beam (with p and s components of equal intensity) into reflected and refracted beams of orthogonal polarizations [Opt. Lett. **31**, 1525 (2006)] and requires a film refractive index $n_1 \geq (\sqrt{2}+1)\sqrt{n_2}$. A monochromatic design that uses a high-index TiO_2 thin film on a low-index MgF_2 substrate at 488 nm wavelength is presented as an example. © 2007 Optical Society of America

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1. INTRODUCTION

In a recent letter¹ it was shown that it is possible to split a monochromatic light beam into two beams of equal power and orthogonal polarizations by reflection and refraction at the planar surface of a dielectric substrate at an angle of incidence below the Brewster angle. This requires a substrate of high refractive index, $n_2 \geq 3 + 2\sqrt{2} = 5.8284$, e.g., PbTe in the IR. It was noted that a quarter-wave layer of high refractive index, which is deposited on a substrate of low refractive index, can also be used to accomplish the same novel beam splitting function. This thin-film coating design is now presented in this paper. For a succinct account of different types of beam splitters, see, e.g., the review by Dobrowolski.²

In Section 2 an analytical design procedure is presented for achieving 50% reflectance of unpolarized (or randomly polarized) light that is obliquely incident on a transparent quarter-wave coating on a transparent substrate. The algorithm developed in Section 2 is applied in Subsection 3.A to the special case of a vanishing substrate ($n_2=1$), i.e., an unsupported quarter-wave pellicle. It is also applied to quarter-wave coatings of varying refractive index n_1 on a (glass or plastic) substrate of refractive index $n_2=1.5$ in Subsection 3.B.

In Section 4 a beam splitter design that uses a TiO_2 high-index quarter-wave coating on a low-index MgF_2 substrate for 488 nm (Ar-ion-laser) light is considered as a specific example. An error analysis shows the effect of $\pm 5^\circ$ shifts in the angle of incidence and of $\pm 5\%$ shifts in wavelength or film thickness on the performance of this beam splitter. Section 5 gives a brief summary of the paper.

2. QUARTER-WAVE LAYERS THAT REFLECT 50% OF UNPOLARIZED LIGHT AT OBLIQUE INCIDENCE: ANALYTICAL TREATMENT

For light reflection in air ($n_0=1$) by a transparent layer of quarter-wave optical thickness (equal to half the film-thickness period) and refractive index n_1 on a transparent substrate of refractive index n_2 , the amplitude reflection coefficients of the p and s polarizations at oblique incidence at an angle φ are real and are given by^{3,4}

$$R_s = \frac{S_0 S_2 - S_1^2}{S_0 S_2 + S_1^2}, \quad (1)$$

$$R_p = \frac{n_1^4 S_0 S_2 - n_2^2 S_1^2}{n_1^4 S_0 S_2 + n_2^2 S_1^2}, \quad (2)$$

where

$$S_i = (n_i^2 - u)^{1/2}, \quad i = 0, 1, 2, \quad (3)$$

$$u = \sin^2 \varphi. \quad (4)$$

All materials are assumed to be optically isotropic and are separated by parallel plane boundaries.

In terms of the quantity

$$P = S_0 S_2 / S_1^2, \quad (5)$$

Eqs. (1) and (2) are rewritten as

$$R_s = \frac{P-1}{P+1}, \quad R_p = \frac{n_1^4 P - n_2^2}{n_1^4 P + n_2^2}. \quad (6)$$

To achieve 50% intensity reflectance for incident unpolarized light, the amplitude reflection coefficients of the p and s polarizations must satisfy the following condition:

$$R_p^2 + R_s^2 = 1. \quad (7)$$

If Eqs. (6) are substituted into Eq. (7), a quartic equation in P is obtained,

$$a_4 P^4 + a_3 P^3 + a_2 P^2 + a_1 P + a_0 = 0, \quad (8)$$

with coefficients given by

$$\begin{aligned} a_4 &= n_1^8, \\ a_3 &= -2n_1^4(n_1^4 + n_2^2), \\ a_2 &= n_1^4(n_1^4 - 12n_2^2) + n_2^4, \\ a_1 &= -2n_2^2(n_1^4 + n_2^2), \\ a_0 &= n_2^4. \end{aligned} \quad (9)$$

For a high-index transparent film on a low-index transparent substrate it is apparent from Eqs. (3)–(5) that P is real and positive and in the range $0 < P < 1$, hence, only roots of Eq. (8) that satisfy this condition are accepted.

From Eqs. (3)–(5) we also obtain

$$P^2 = (1-u)(n_2^2 - u)/(n_1^2 - u)^2. \quad (10)$$

Equation (10) can be rewritten as a quadratic equation in u :

$$b_2 u^2 + b_1 u + b_0 = 0, \quad (11)$$

with coefficients given by

$$\begin{aligned} b_2 &= P^2 - 1, \\ b_1 &= (n_2^2 + 1) - 2n_1^2 P^2, \\ b_0 &= n_1^4 P^2 - n_2^2. \end{aligned} \quad (12)$$

The only acceptable roots of Eq. (11) are those for which $0 \leq u < 1$. Finally, the angle of incidence φ at which the 50% reflectance for incident unpolarized light is achieved, for given n_1 and n_2 , is given by

$$\varphi = \arcsin(u^{1/2}). \quad (13)$$

3. QUARTER-WAVE LAYERS THAT REFLECT 50% OF UNPOLARIZED LIGHT AT OBLIQUE INCIDENCE: NUMERICAL RESULTS

A. Case of a Vanishing Substrate, $n_2=1$

The algorithm of Section 2 is applied to determine the constraint between n_1 and φ such that Eq. (7) is satisfied when $n_2=1$ (i.e., for an unsupported quarter-wave layer

or pellicle), by assigning values to n_1 between 1 and 6 in steps of 0.01 and solving for the corresponding values of φ .

Figure 1 shows the two solution branches that we obtain: a high-angle branch (HAB), or solution 1, which is represented by the continuous line, and a low-angle branch (LAB), or solution 2, which is represented by the dashed line. The dash-dot curve in the middle represents the polarizing (Brewster) angle, $\varphi_p = \varphi_B = \tan^{-1}n_1$.

Figures 2 and 3 show the real amplitude reflection coefficients R_p and R_s that are associated with the HAB and LAB, respectively, as functions of the film refractive index n_1 . R_p and R_s are both negative for the HAB, Fig. 2, and have opposite signs ($R_p > 0, R_s < 0$) for the LAB, Fig. 3.

Figure 4 is a plot of R_p versus R_s and shows two separate arcs of the unit circle, Eq. (7), in the third and fourth quadrants, respectively. Points A and B correspond to the similarly marked points in Fig. 1.

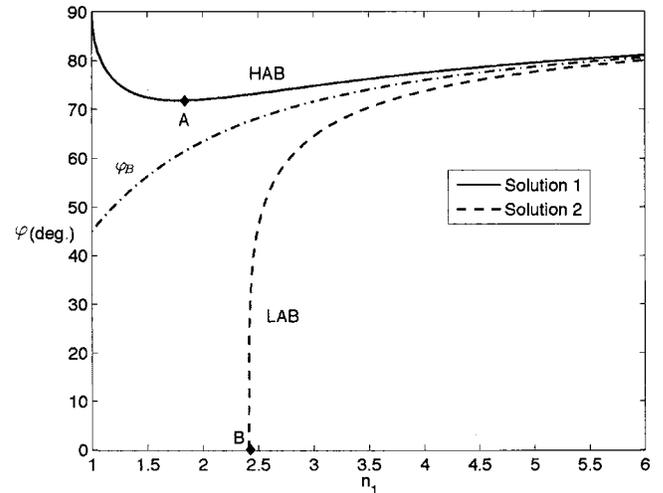


Fig. 1. Angle of incidence φ versus film refractive index n_1 such that 50% of incident unpolarized light is reflected by a quarter-wave pellicle. The continuous and dashed curves represent two independent solution branches. The middle curve gives the Brewster angle as a function of n_1 , $\varphi_B = \tan^{-1}n_1$.

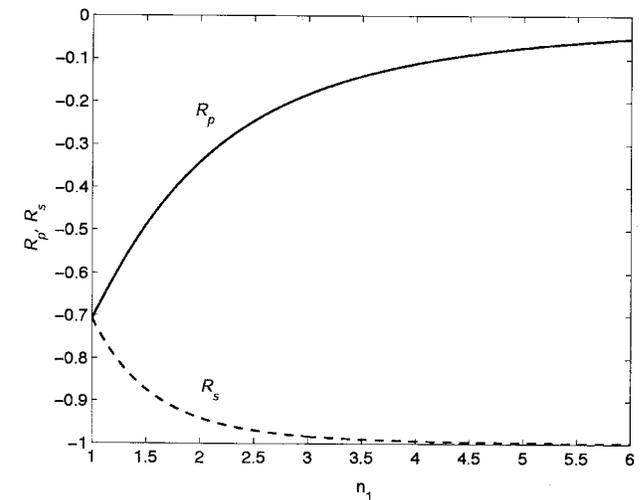


Fig. 2. Amplitude reflection coefficients R_p and R_s that are associated with the HAB of Fig. 1 plotted as functions of the film refractive index n_1 .

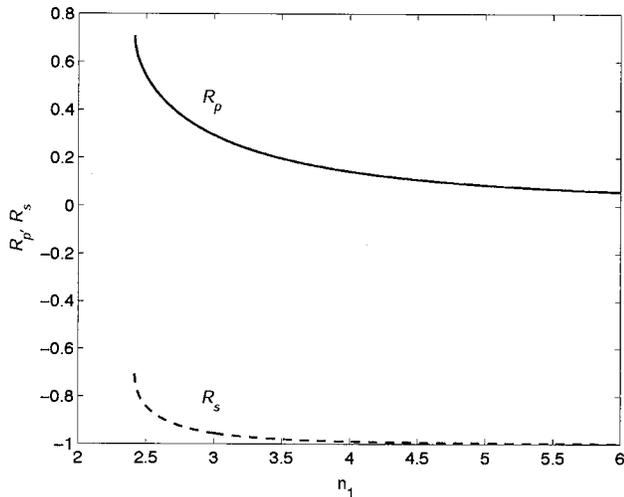


Fig. 3. Amplitude reflection coefficients R_p and R_s that are associated with the LAB of Fig. 1 plotted as functions of the film refractive index n_1 .

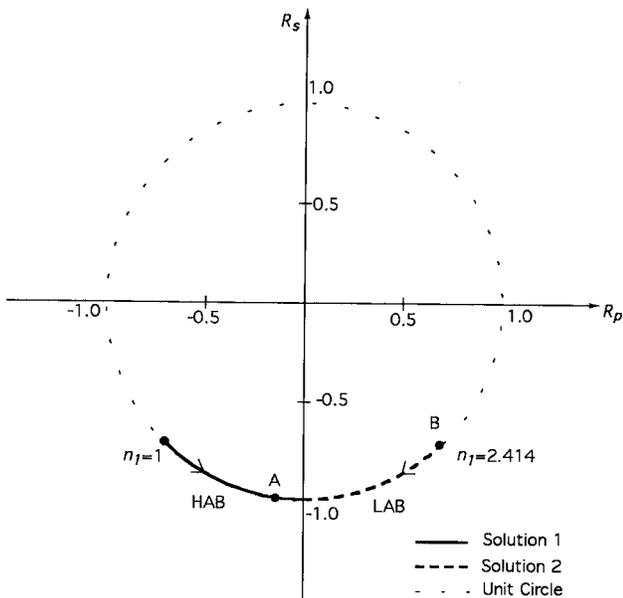


Fig. 4. Plot of R_p versus R_s associated with the high- and low-angle branches (Figs. 2 and 3) yields two arcs of the unit circle in the third and fourth quadrants, respectively. The arrows indicate the direction of increasing n_1 , and points A and B correspond to the similarly marked points in Fig. 1.

Figures 5 and 6 show the intensity reflectances R_p^2 and R_s^2 for the HAB and LAB, respectively, as functions of n_1 . The average value of R_p^2 and R_s^2 is constant ($=0.5$), independent of n_1 as is required by Eq. (7).

An interesting feature of Fig. 1 is that φ of the HAB reaches a minimum, $\varphi_{\min}=71.8225^\circ$, at point A where $n_1=1.78$. Therefore, a pellicle of 1.78 refractive index and quarter-wave optical thickness (or an odd integral multiple thereof) functions as a 50%-50% BS for incident unpolarized light at $\approx 72^\circ$ angle of incidence. The starting point B of the LAB in Fig. 1 is at $\varphi=0$ and corresponds to $n_1=\sqrt{2}+1=2.414$. Film materials that have this refractive index include diamond in the visible and ZnSe in the IR.^{5,6}

B. General Case of $n_2 > 1$

For an arbitrary substrate of refractive index n_2 , the starting value of n_1 for the LAB is obtained by noting that at $\varphi=0$, $R_p^2=R_s^2=1/2$. It follows from Eqs. (6) that

$$R_s = (P - 1)/(P + 1) = -1/\sqrt{2}, \tag{14}$$

$$P = (\sqrt{2} - 1)^2. \tag{15}$$

If $u=0$ is substituted in Eq. (10), we get

$$n_1 = (n_2/P)^{1/2}. \tag{16}$$

When Eqs. (15) and (16) are combined, we obtain

$$n_1 = (\sqrt{2} + 1)\sqrt{n_2}. \tag{17}$$

As a specific example we take quarter-wave layers on a (glass or plastic) substrate with $n_2=1.5$. Figure 7 shows the two solution branches that we obtain using the algorithm of Sec. 2. When compared with Fig. 1 (for $n_2=1$), the HAB in Fig. 7 is displaced upward to higher angles,

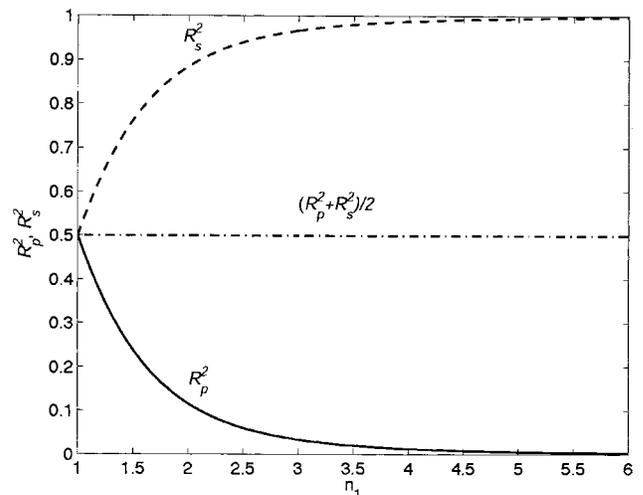


Fig. 5. Intensity reflectances R_p^2 and R_s^2 that are associated with the HAB of Fig. 1 plotted as functions of the film refractive index n_1 . The average reflectance, $(R_p^2+R_s^2)/2$, is constant at 0.5.

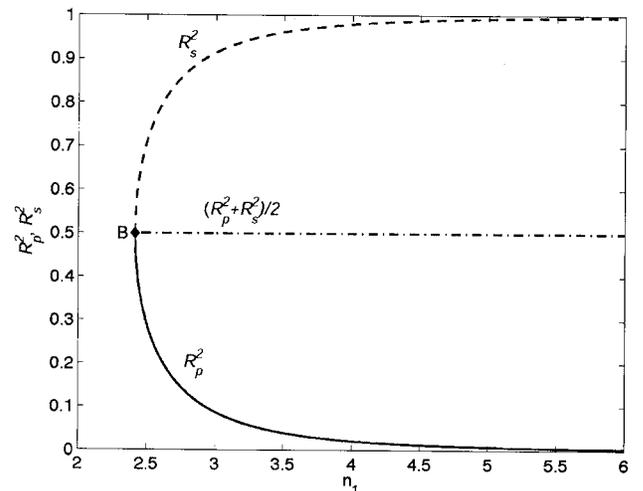


Fig. 6. Intensity reflectances R_p^2 and R_s^2 that are associated with the LAB of Fig. 1 plotted as functions of the film refractive index n_1 . The average reflectance, $(R_p^2+R_s^2)/2$, is constant at 0.5.

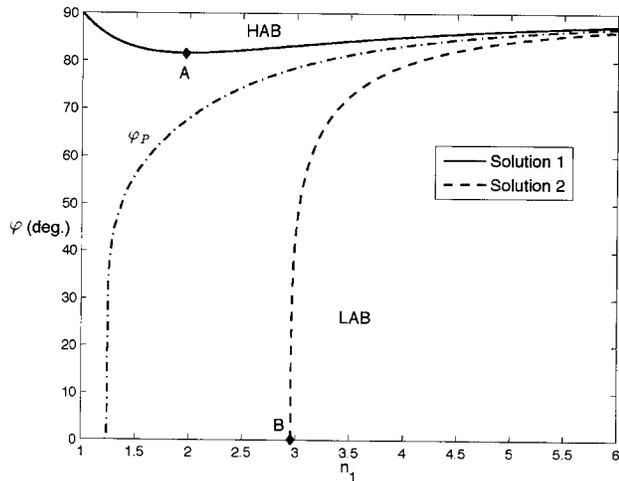


Fig. 7. Angle of incidence φ versus the refractive index n_1 of a transparent quarter-wave coating on a transparent substrate ($n_2=1.5$) that reflects 50% of incident unpolarized light. The continuous and dashed curves represent two distinct solution branches. The middle dash-dot curve gives the polarizing angle φ_p as a function of n_1 as obtained from Eqs. (19) and (20).

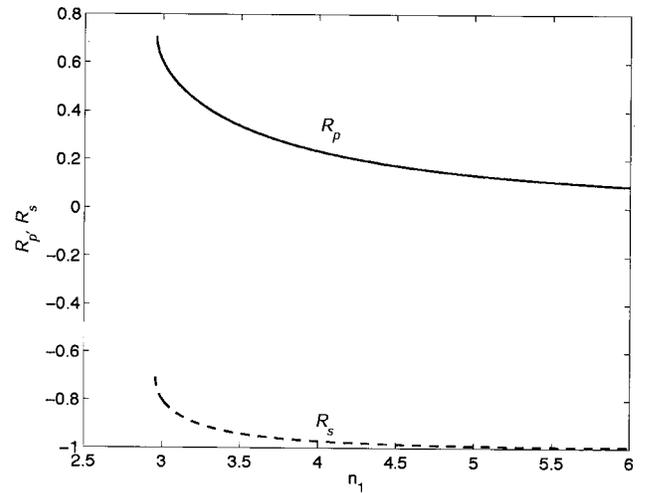


Fig. 9. Amplitude reflection coefficients R_p and R_s that are associated with the LAB of Fig. 7 plotted as functions of the film refractive index n_1 .

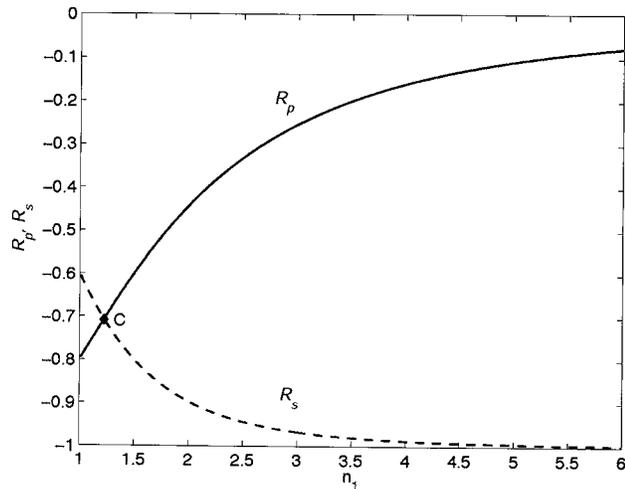


Fig. 8. Amplitude reflection coefficients R_p and R_s that are associated with the HAB of Fig. 7 plotted as functions of the film refractive index n_1 . The significance of the point of intersection C is discussed in the text.

and its minimum at point A now occurs at $\varphi_{\min}=81.6208^\circ$, $n_1=1.98$. The starting value $n_1=2.9568$ of the LAB at point B is predicted by Eq. (17).

In Fig. 7 the middle (dash-dot) curve gives the polarizing angle φ_p of the film-substrate system at which $R_p=0$ as a function of the film refractive index n_1 with $n_2=1.5$. To calculate the polarizing angle φ_p we note from Eqs. (6) that $R_p=0$ if

$$P = n_2^2/n_1^4. \quad (18)$$

Substitution of P from Eq. (18) into Eq. (10) gives a quadratic equation,

$$c_2u^2 + c_1u + c_0 = 0, \quad (19)$$

in $u = \sin^2 \varphi_p$ with coefficients given by

$$c_2 = n_1^8 - n_2^8,$$

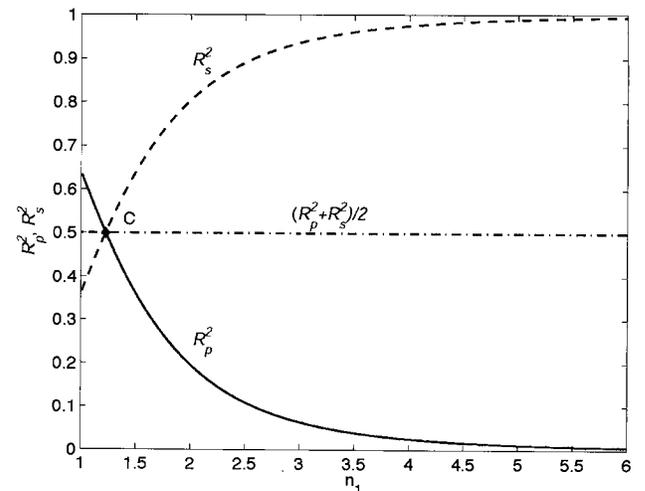


Fig. 10. Intensity reflectances R_p^2 and R_s^2 that are associated with the HAB of Fig. 7 plotted as functions of the film refractive index n_1 . The average reflectance, $(R_p^2 + R_s^2)/2$, is constant at 0.5.

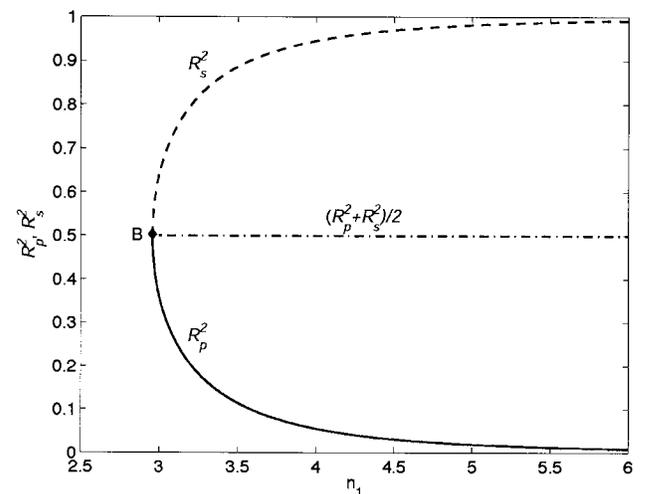


Fig. 11. Intensity reflectances R_p^2 and R_s^2 that are associated with the LAB of Fig. 7 plotted as functions of the film refractive index n_1 . The average reflectance, $(R_p^2 + R_s^2)/2$, is constant at 0.5.

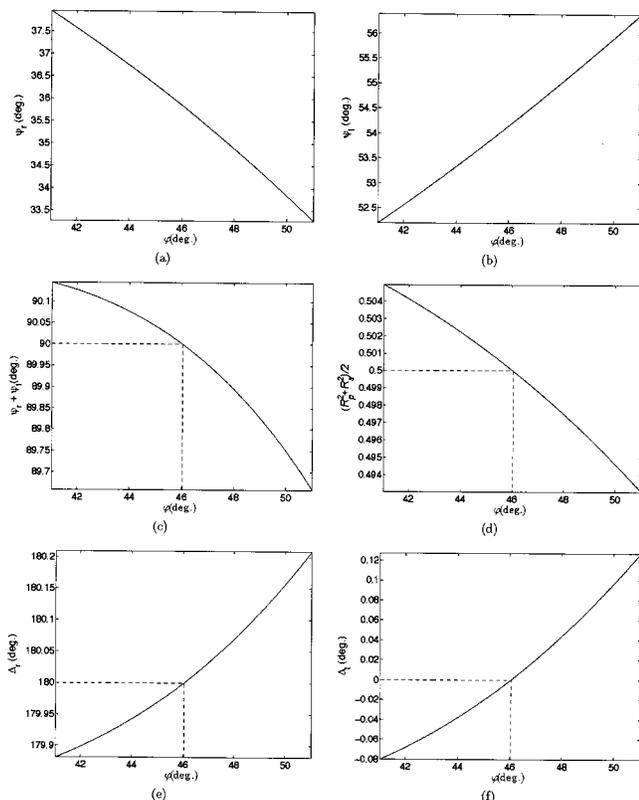


Fig. 12. Average reflectance $(R_p^2 + R_s^2)/2$ and ellipsometric parameters $\psi_r, \psi_t, \Delta_r, \Delta_t$ as functions of the angle of incidence φ for a beam splitter that consists of a quarter-wave (43.51 nm) thin film of TiO_2 on a MgF_2 substrate at 488 nm wavelength. The angle of incidence is varied by $\pm 5^\circ$ around the design angle $\varphi = 46.04^\circ$.

$$c_1 = 2n_1^2 n_2^4 - n_1^8 (n_2^2 + 1),$$

$$c_0 = n_1^4 n_2^2 (n_1^4 - n_2^2). \quad (20)$$

Figures 8 and 9 show the amplitude reflection coefficients R_p and R_s as functions of n_1 for the HAB and LAB of Fig. 7, respectively. Again, R_p and R_s are both negative for the HAB, Fig. 8, and have opposite signs ($R_p > 0, R_s < 0$) for the LAB, Fig. 9.

The point of intersection C of the curves of R_p and R_s in Fig. 8 is located at

$$R_p = R_s = -1/\sqrt{2} = -0.7071, \quad n_1 = \sqrt{n_2} = 1.2247. \quad (21)$$

This corresponds to a polarization-preserving (or polarization-independent) 50%–50% beam splitter in both reflection and transmission.³

The associated intensity reflections R_p^2 and R_s^2 as functions of n_1 are shown in Figs. 10 and 11 for the HAB and LAB, respectively. Again, the average of R_p^2 and R_s^2 is constant (=0.5) independent of n_1 as required by Eq. (7).

4. BEAM SPLITTER FOR PRODUCING REFLECTED AND TRANSMITTED BEAMS OF EQUAL POWER AND ORTHOGONAL POLARIZATIONS IN THE VISIBLE

Consider a high-index ($n_1 = 2.895$) thin film of TiO_2 on a low-index ($n_2 = 1.386$) MgF_2 substrate at wavelength λ

=488 nm. (Because we assume isotropic phases, the refractive indices are taken as averages of the ordinary and extraordinary indices n_o and n_e published in Ref. 6.) By use of the algorithm of Section 2, the operating angle of incidence $\varphi = 46.037^\circ$ is obtained. The film-thickness period D_1 and quarter-wave metric film thickness d_1 are determined by

$$D_1 = (\lambda/2)(n_1^2 - \sin^2 \varphi)^{-1/2},$$

$$d_1 = D_1/2, \quad (22)$$

which give $D_1 = 87.02$ nm and $d_1 = 43.51$ nm. Under the ideal operating conditions, the average reflectance for incident unpolarized light is 1/2 and the reflection and transmission ellipsometric parameters⁴ satisfy the following relations:^{1,7}

$$\Delta_r = \pi, \Delta_t = 0,$$

$$\psi_r + \psi_t = \pi/2. \quad (23)$$

To couple the refracted light out of the substrate into air, a prismatic substrate with prism angle equal to the angle of refraction (31.287°) can be used, so that light leaves the prism normal to its exit face which is antireflection coated.⁸

When the wavelength of light and metric film thickness are kept constant, and the angle of incidence is varied by up to $\pm 5^\circ$, the average reflectance and ellipsometric pa-

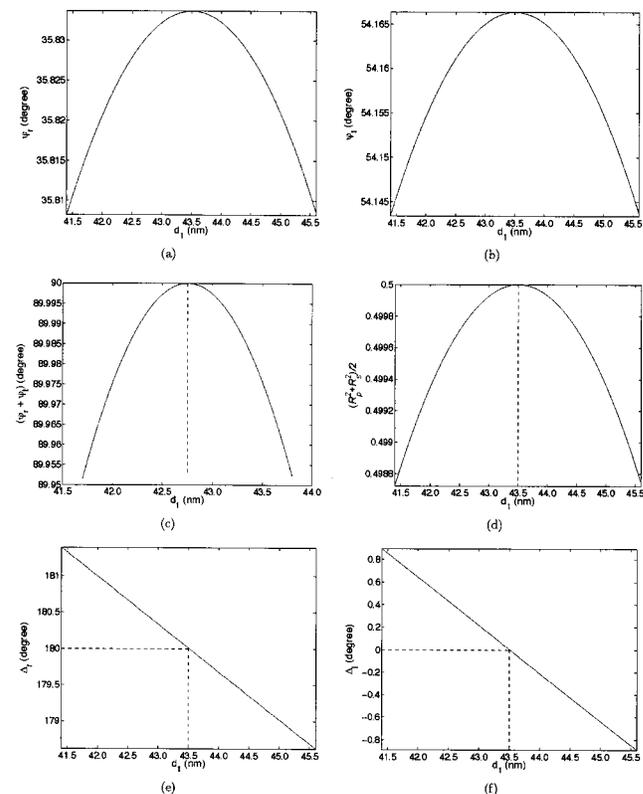


Fig. 13. Average reflectance $(R_p^2 + R_s^2)/2$ and ellipsometric parameters $\psi_r, \psi_t, \Delta_r, \Delta_t$ of a beam splitter as functions of the thickness d_1 of a TiO_2 film on a MgF_2 substrate at $\lambda = 488$ nm and $\varphi = 46.04^\circ$. The film thickness is varied by $\pm 5\%$ around the design value of 43.5 nm.

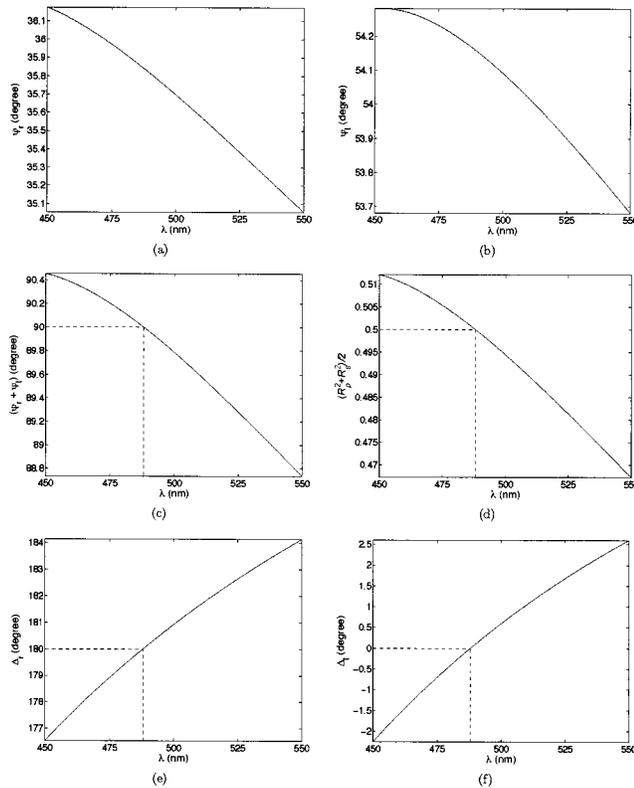


Fig. 14. Average reflectance $(R_p^2 + R_s^2)/2$ and ellipsometric parameters ψ_r , ψ_i , Δ_r , Δ_i as functions of the wavelength λ (in nanometers) for a beam splitter that consists of a 43.51 nm TiO_2 thin film on a MgF_2 substrate at an angle of incidence of $\varphi = 46.04^\circ$. The wavelength λ is changed by $\pm 5\%$ around 488 nm.

rameters change in the manner shown in Fig. 12. From Fig. 12, it is apparent that the ideal conditions of Eqs. (23) are nearly maintained; hence, this novel beam splitter is tolerant to small angle-of-incidence errors.

If the wavelength of light and angle of incidence are fixed, but the film thickness is varied by $\pm 5\%$ around its design value ($d_1 = 43.51$ nm), the average reflectance and ellipsometric parameters change in the manner shown in Fig. 13. Again, the ideal conditions of Eqs. (23) are approximately satisfied, to within a small error, and the coating design is not overly sensitive to small film-thickness errors around the quarter-wave thickness.

Finally, when the angle of incidence and metric thickness of the coating are kept constant, and the wavelength

is shifted by $\pm 5\%$ around $\lambda = 488$ nm, the average reflectance and ellipsometric parameters change in the manner shown in Fig. 14. In Fig. 14, the deviations from the ideal conditions of Eqs. (23) are more pronounced; hence, the beam splitter is not considered achromatic.

5. SUMMARY

A transparent layer of quarter-wave optical thickness, a pellicle or a coating deposited on a transparent substrate, reflects 50% of incident unpolarized light under certain conditions that are completely determined. Two distinct solution branches are obtained that correspond to incidence below and above the p -suppressing polarizing angle of the film-substrate system. Operation below the polarizing angle makes possible a novel thin-film beam splitter¹ that divides incident totally polarized light into reflected and refracted beams of equal power (50% each) and orthogonal polarizations. The performance of one such device that uses a high-index TiO_2 quarter-wave coating on a low-index MgF_2 substrate at 488 nm wavelength is presented.

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